

Emissivity Measurements on Shape Memory Alloys

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Abstract

In many applications of Shape Memory Alloys, the most appropriate means and sometimes the only one to measure temperature is thermography, which relies on knowing the emissivity of the material. This paper aims to evaluate experimentally the behavior of the emissivity of a NiTi shape memory alloy as a function of temperature in order to investigate how the emissivity behaves during phase transformation and how it may impact the measured temperature. The emissivity of a Ni₅₅Ti₄₅ shape memory alloy was measured by means of active and passive infrared sensors in a spectral band of 8-14 μm in the temperature range from 25 to 100°C under a vacuum of 3×10^{-3} Pa. The results show that the emissivity of the Ni₅₅Ti₄₅ shape memory alloy varied approximately 9% between fully martensitic and fully austenitic states.

1. Introduction

In the last decades, the use of intelligent materials has been intensified in a variety of mechanical systems [1-2]. Among those systems, the *shape memory alloys* (SMAs) show a great potential in the shape memory effect, especially equiatomic NiTi alloys, which can deform itself up to 8% and restore itself to its original shape, and then repeat the same operation a large number of times [3]. In certain applications such as the actuators that involve important deformations, loads and low frequencies, the SMAs have a significant potential in relation to the traditional actuators due to their high proportion of energy and silent performance [4]. Furthermore, those alloys show a thermomechanical behaviour highly dependent on temperature as well as four characteristic temperatures: M_s (*Martensite start*) initial temperature to form martensite; M_f (*Martensite finish*) final temperature to form martensite; A_s (*Austenite start*) initial temperature to form austenite; A_f (*Austenite finish*) final temperature to form austenite [5]. Notably, in both thermomechanical behaviour and many other applications measuring and quantifying the temperature of the shape memory elements is fundamental [6].

Depending on the SMAs' application, the conventional methods for measuring temperature such as thermocouples sometimes do not provide adequate results, mainly due to the effects of contact resistance loss. Alternatively, thermography is a technique that has been utilised with many advantages when contrasted to conventional methods for measuring temperature inasmuch as it allows superficial temperature variations to be determined from a distance and contactless [7-8].

Even though thermography is the most adequate technique to quantify temperature variations remotely, the results obtained through that method are highly dependent on various physical factors related to radiant properties of material, mainly thermal emissivity [9]. Therefore, knowing the emissivity is necessary. The aim of this project is to analyse the variation caused by temperature in total hemispheric emissivity of a shape memory alloy (Ni₅₅Ti₄₅). The first part of the study, which presented preliminary experimental results, has already been published [10]. The following work is a correspondent investigation.

2. Materials and Methods

2.1. Samples

For this study, we utilized two samples of NiTi alloy with nominal composition of 55% in Ni weight supplied by Nemesis Technology. The samples were produced and prepared in accordance with JIS A 1423 (Simplified test method for emissivity by Infrared Radio Meter) with cylindrical shape [11]. We opted for that format due to difficulties in the machining process, among which its hardness (66.84 HRA) stands out, as well as the existence of data showing that the emissivity of the cylinder-shaped SMA was obtained through thermography in similar conditions [12]. Moreover, cutting the samples can create superficial tensions that can eventually result in martensitic transformations of the material [13-14]. The cuts A₁ and A₂ were made with a precision cut-off machine (*Struers Secotom-15*) in low-speed conditions (0.005 mm/s) with around 10 mm diameter and 100 mm length, being continuously cooled by a cooling fluid. Afterwards, the samples received additional care such as ultrasonic cleaning bath during 20 minutes in a (CH₃)₂CO solution.

Table 1 exposes the phase transformation temperatures of the alloy measured by a DSC-8500 and Pyris® Perkin Elmer software with heating and cooling rates controlled at 20°C/minute under continuous nitrogen gas flow. The

general alloy composition, “as-received” obtained by EDX-720, is presented on Table 2. A small quantity of Al, Fe, Ca and Si is present.

Table 1. Phase transformation temperatures.

Sample Ni ₅₅ Ti ₄₅	A _s (°C)	A _f (°C)	M _s (°C)	M _f (°C)
A ₁	57.01	74.80	41.41	15.62
A ₂	57.00	75.10	41.20	15.30

Table 2. Chemical composition of the NiTi alloy “as-received”.

Chemical content	Ni	Ti	Al	Fe	Ca	Si
% wt.	55.408	43.888	0.295	0.157	0.143	0.109

The calorimetric analysis showed that each sample is partially martensitic at ambient temperature. Thus, roughness was measured at stable ambient temperature. Arithmetic average surface roughness (R_a), stands for the root mean square (R_q) and for the maximum peak to valley distance in profile (R_t) were measured for each sample. Table 3 shows the roughness properties for the analyses carried out at ambient temperature.

Table 3. Sample roughness at ambient temperature.

Sample 25° C	R _a	R _t	R _q
A ₁	0.761	5.41	0.968
A ₂	0.760	5.36	0.960

At ambient temperature, around 25°C, samples A₁ and A₂ showed very similar superficial characteristics.

2.2. Emissivity measurements

The emissivity of samples A₁ and A₂ as a function of temperature was measured by an experimental apparatus developed for that purpose. Figure 1 shows the device, built in accordance with JIS A 1423 (Simplified test method for emissivity by Infrared Radio Meter) and ASTM E1933-99 (Standard Test Methods for Measuring and Compensating for Emissivity Using Infrared Imaging Radiometers) [15]. It carries out measurements through direct or indirect calorimetric and radiometric methods for temperatures that vary from 25 to 100°C, and for wavelengths between 8 and 14 μm under vacuum. In order to measure total hemispheric emissivity (integrated in 180° from θ=-90° to θ=90°), the apparatus uses an active source and an infrared detector (IR) connected to the controlling block through optical fibre cables. When the IR source is on, it has a commutation frequency that varies from 0.5 to 20 Hz. Then a microcontroller stores the data collected from its analog-to-digital converter, called ADCI and ADCII, in E²PROM memory. After that, the processor calculates the difference between the IR radiation in ADCI/ADCII and the radiation received during IR transmitter variation – converters called ADCIII and ADCIV. Therefore, emission rate is obtained through the equation (1).

$$\text{Epsout} = 1 - \frac{(\text{ADCIII} - \text{ADCIV})}{(\text{ADCI} - \text{ADCII})} * fm - fR \quad (1)$$

where fm is the reflection index and fR is the hemispheric correction factor.

The configuration applied permitted us to obtain the emissivity dependence angle (angle between emission direction and normal on the sample surface) in relation to its inclination. The sample was placed on a support and a Joule effect heater device was used for heating it. During the testing, when the sample emissivity value changes, the new value is calculated and transferred through an analogical output 0-10 volts to a software. Thus, the signal is conditioned to a range where: 0 V ϵ = 0.100 and 10 V ϵ = 1.00.

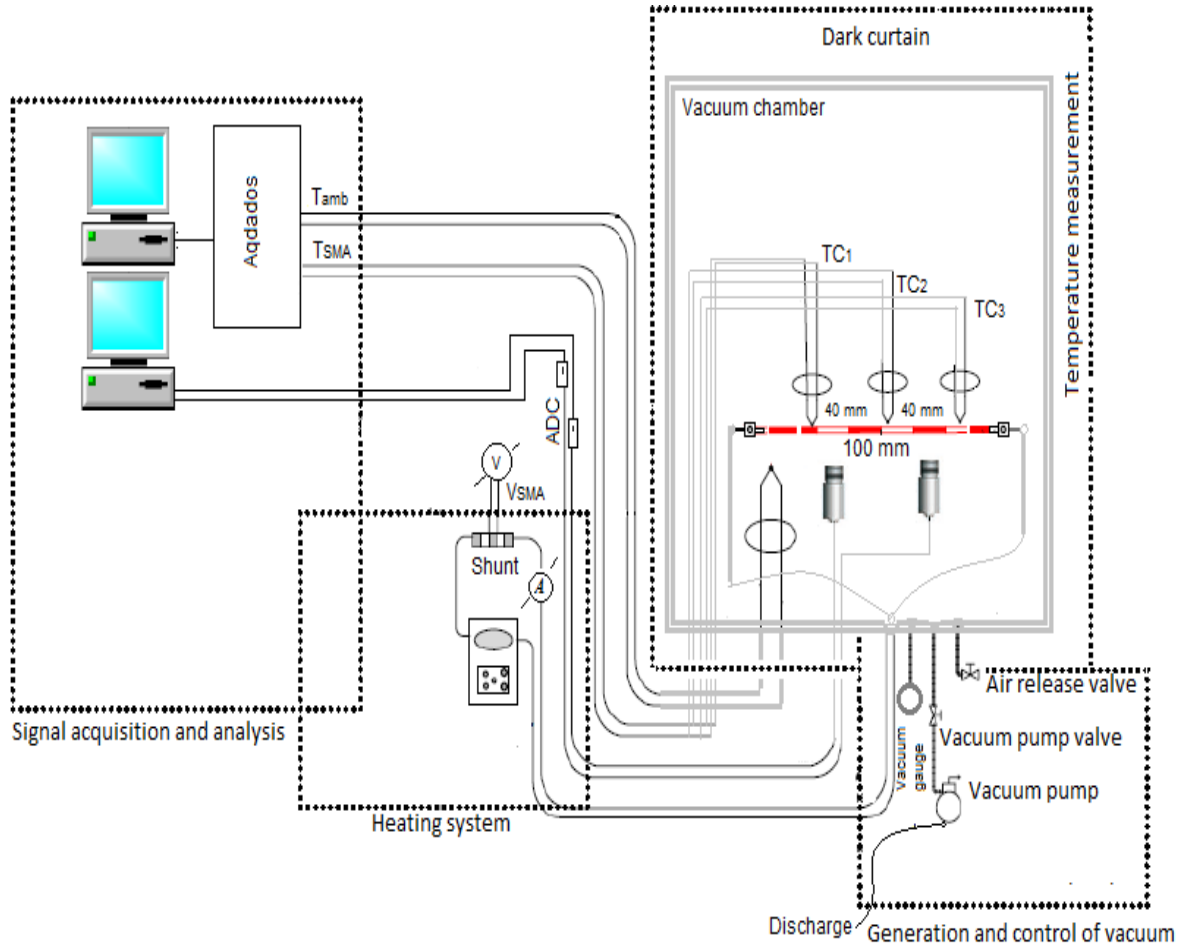


Fig. 1. Schematic of apparatus developed for measuring emissivity.

Three type K thermocouples welded to the sample surface, placed away from the visualisation area of the IR sensors, measured temperature. The average of temperatures was called T_1 . Three PTC sensors were utilised to obtain the temperature inside the vacuum bell jar. The average of temperatures here was named T_2 . Pressure inside the vacuum bell jar is kept at 3×10^{-3} Pa during the experiments. Under these conditions, assuming that the heat loss throughout the sample due to conduction and convection are insignificant when compared to the loss caused by radiation, the emissivity is [16]:

$$\varepsilon_1 = \frac{Q}{\sigma A_1 (T_1^4 - T_2^4)} \quad (2)$$

where ε_1 is the total hemispheric emissivity of the sample in T_1 ; $Q = IV$ (W) is the electric power obtained in the sample; σ is the Stefan-Boltzmann constant; A_1 is the surface area where the tests were carried out and T_1 and T_2 are the temperatures of the given surface area and the wall under vacuum respectively.

The total hemispheric emissivity in Equation (2) is a function of area A_1 , temperatures T_1 and T_2 and electric power calculated through the voltage drop, V , and current I [17]. The function is:

$$\varepsilon = F(T_1, T_2, A_1, Q) \quad (3)$$

The total hemispheric emissivity uncertainty is measured by:

$$\frac{\Delta \varepsilon}{\varepsilon} = \left[\left(\frac{\partial \ln F}{\partial S} \right)^2 (\Delta S)^2 + \left(\frac{\partial \ln F}{\partial T_1} \right)^2 (\Delta T_1)^2 + \left(\frac{\partial \ln F}{\partial T_2} \right)^2 (\Delta T_2)^2 + \left(\frac{\partial \ln F}{\partial U} \right)^2 (\Delta U)^2 + \left(\frac{\partial \ln F}{\partial I} \right)^2 (\Delta I)^2 \right]^{1/2} \quad (4)$$

3. Experimental Results and Discussion

The total hemispheric emissivity of the samples A₁ and A₂ was measured from 25 to 100°C approximately using the apparatus describe in section 2.2 and starting in total martensitic state. The curves and analyses were carried out at MATLAB [18]. Figure 2 shows total hemispheric emissivity curves as a function of temperature. Data reveal the average of direct radiometric measurements for the samples A₁ (Figure 2a) and A₂ (Figure 2b). The results in Figure 2 show that the total hemispheric emissivity of the samples A₁ and A₂ during heating is almost constant for temperatures varying from 26 to 65°C and decreases between temperatures A_s and A_f with temperature rise. The average emissivity variation of the sample A₁ between temperatures A_s and A_f was around 19% in direct measurements (Figure 2a) and also 19% in indirect measurements (Figure 3a). For the sample A₂, the average emissivity variation between temperatures A_s and A_f was around 15.2% in direct measurements (Figure 2b) and 6.1% in indirect measurements (Figure 3b).

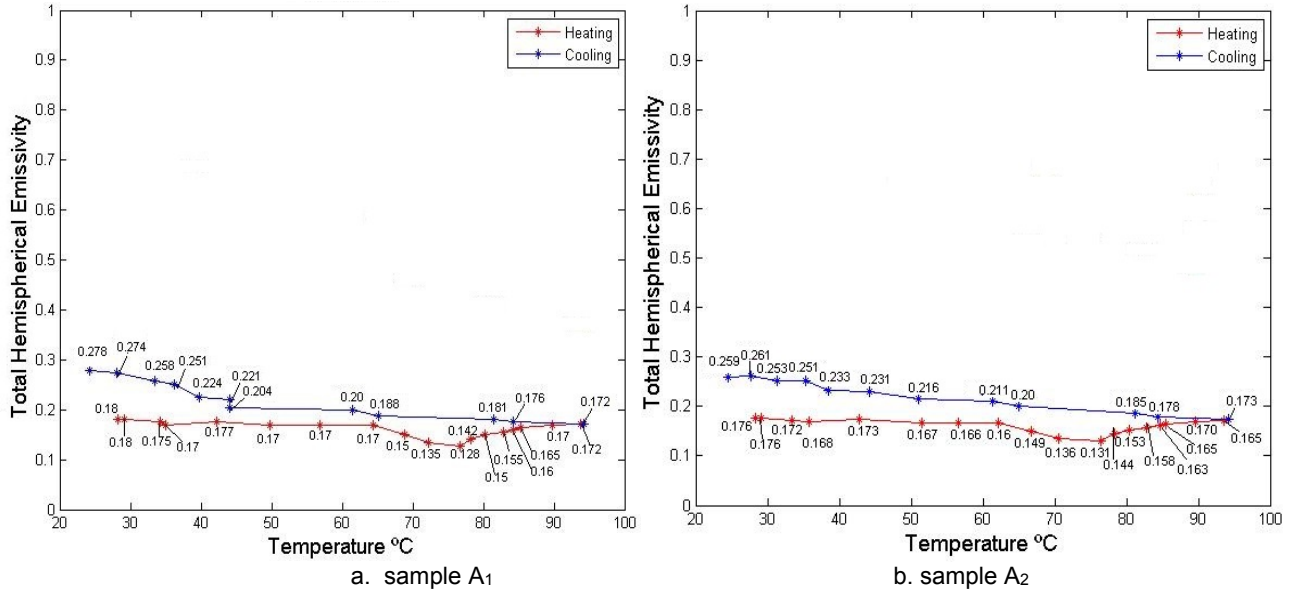


Fig. 2. Total hemispheric emissivity curves registered in direct measurements.

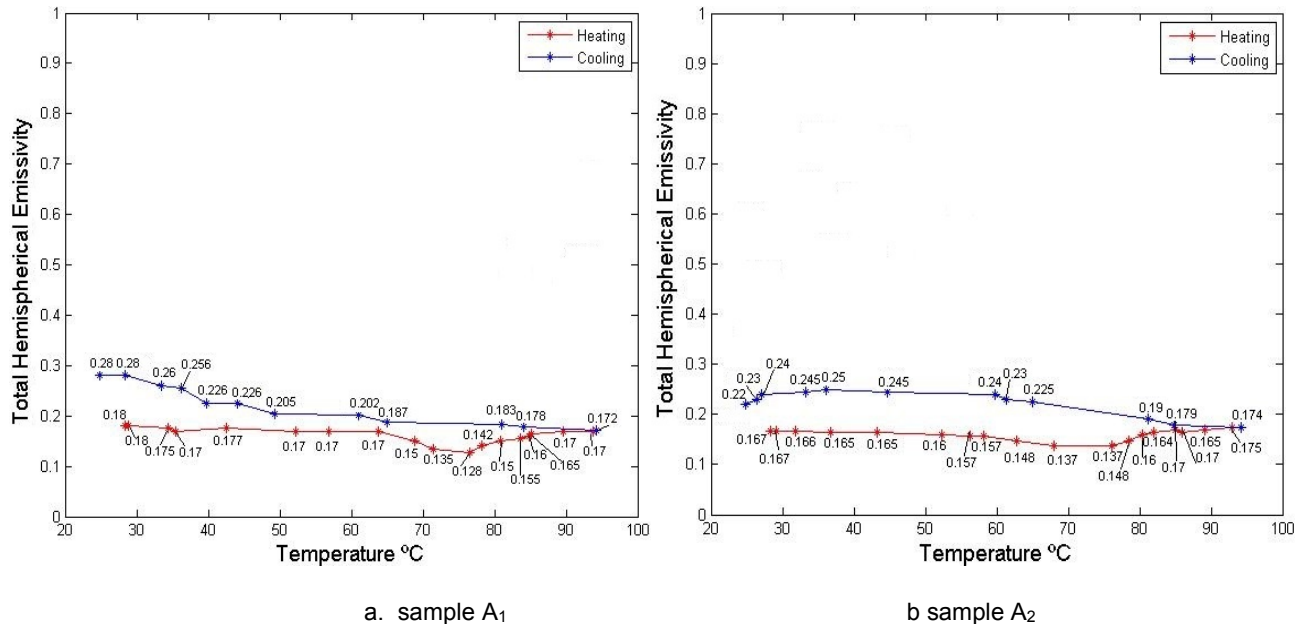


Fig. 3. Total hemispheric emissivity curves registered in indirect measurements.

During the heating phase, emissivity decreased from 0.171 in martensitic phase to 0.155 in austenitic phase (Figure 2). Hemispheric emissivity variation can be explained theoretically and experimentally through the average roughness variation of the alloy – which is 0.760 μm in martensitic phase and 0.821 μm in austenitic phase [19]. At

temperature A_s emissivity starts to decrease due to the reduction of martensitic relief and emissivity directions since each surface protuberance caused by martensitic variant has its own normal direction [20]. At temperature A_f martensitic relieves tend to completely disappear, minimizing emissivity. At temperature over A_f emissivity increases again probably as a result of temperature rise. During the cooling phase, emissivity behaviour is different, that is, emissivity values increase and starting with M_s those values tend to increase more intensely.

Figure 4 shows the two manners of NiTi surface reflection. On the one hand, in martensitic state the sample surface behaves like a diffuse emitter and reflector, that is, there is a larger quantity of directions for emission to be radiated due to more significant roughness (Figure 4b). On the other hand, when the phase is totally austenitic (Figure 4a) reflection is regular, keeping more parallelism during the incidence of rays and minimizing directions.

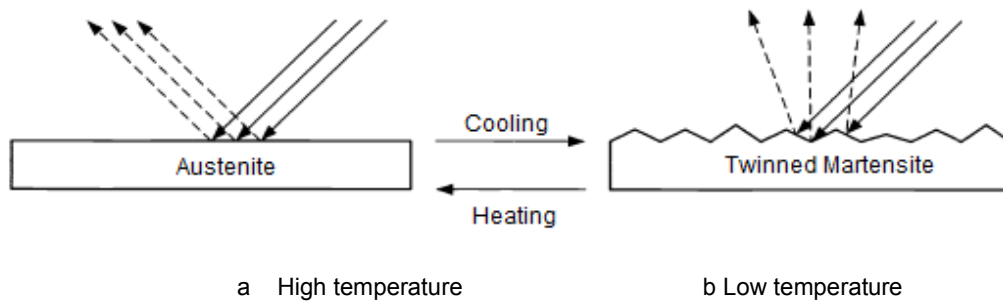


Fig. 4. Representation of NiTi surface reflection [19].

Figure 5 shows emissivity curves superposition as a function of temperature during the heating phase (Figures 2 and 3) and Figure 6 shows the total hemispherical emissivity curve obtained through calorimetric measurements. The data of the Figure 6 curve refer to the average of values in sample A₁ and were calculated by the Equation (2).

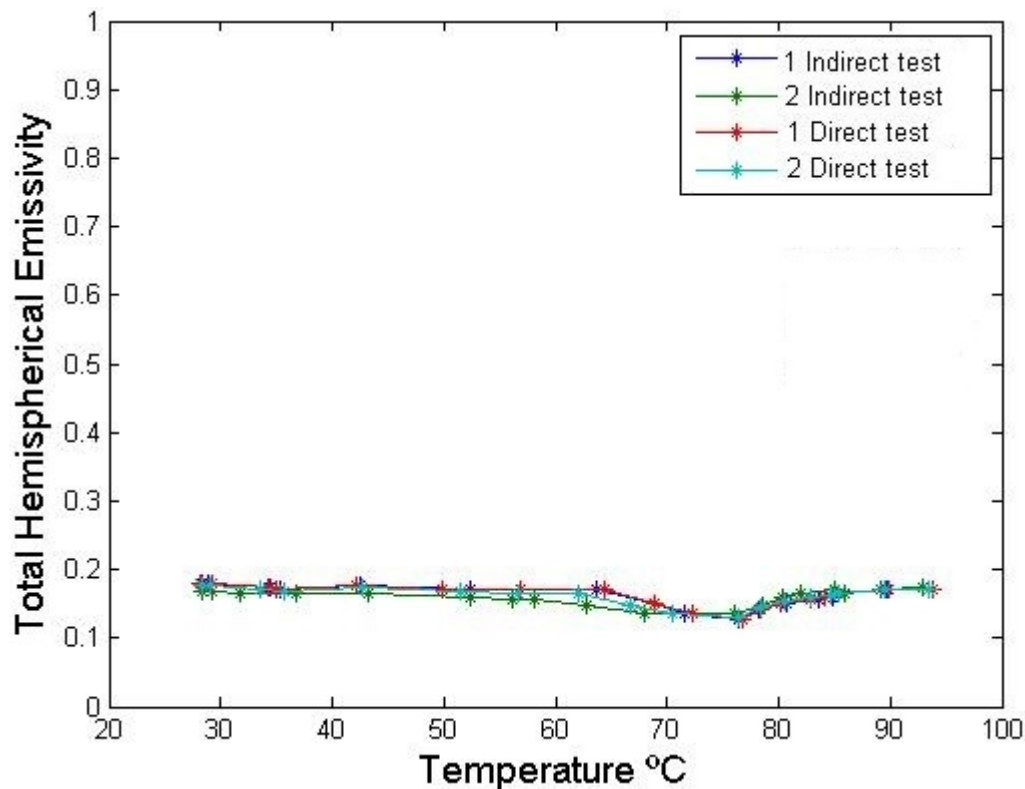


Fig. 5. Emissivity curves superposition versus temperature during heating phase (Figures 2 and 3).

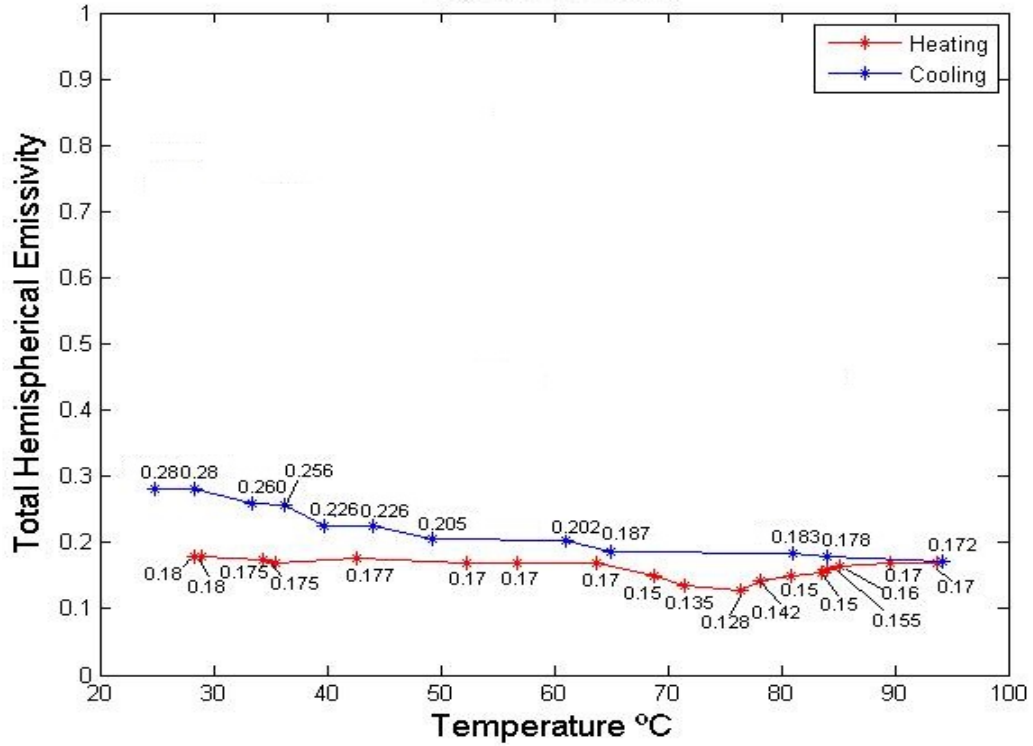


Fig. 6. Total hemispheric emissivity curve versus temperature of the sample A1.

For the temperature range analysed, data reveal that the most expressive total hemispheric emissivity variation happens in phase change. Emissivity varies approximately 24% from M_s at ambient temperature (24°C). There is a correlation between emissivity variation and the error produced when temperature was measured if the emissivity is not corrected. That error is calculated by:

$$T^4 = \frac{E(t)}{\epsilon\sigma}, \tag{5}$$

Replacing temperature and emissivity in the equation, we have:

$$\frac{T_1^4}{T_2^4} = \frac{\frac{E(t)}{\epsilon_1\sigma}}{\frac{E(t)}{\epsilon_2\sigma}} \tag{6}$$

where T_1 is the average temperature measured before adjusting emissivity; T_2 is the real temperature on the surface; ϵ_1 is the emissivity measured before the adjustment; and ϵ_2 is the real emissivity of the surface. The solution of the equation is:

$$T_2 = \sqrt[4]{\frac{\epsilon_1}{\epsilon_2}} T_1 \tag{7}$$

For emissivity experimental data in this work, we have:

$$T_2 = 1,0377T_1. \tag{8}$$

Therefore, temperature T_1 at the final temperature for austenite formation (A_f) read as 78.2°C is then corrected to 81.15°C with emissivity adjustment ϵ_2 . Thus, 15% of emissivity variation corresponds to a reading error of approximately 3 °C without emissivity correction.

4. Conclusions

The total hemispheric emissivity of two SMA alloys was experimentally analysed as a function of temperature through diverse methods. We could understand that global average hemispheric emissivity showed variations of around 9% between martensitic and totally austenitic states. Preliminary analysis of superficial characteristics of the alloy reveals that the variation occurred mainly as a result of surface relief associated with martensitic phase transformation.

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